## Electronic supplementary materials

# Geometrical transition properties of vortex cavitation and associated flow-choking characteristics in poppet valves 

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## S1 Complete form of the momentum equations for 2D axisymmetric geometries

Under the axisymmetric assumption, there are no circumferential gradients in the flow, but circumferential velocities are permitted.

$$
\begin{array}{r}
\frac{\partial}{\partial t}\left(\rho_{\mathrm{m}} u\right)+\frac{1}{r} \frac{\partial}{\partial z}\left(r \rho_{\mathrm{m}} u^{2}\right)+\frac{1}{r} \frac{\partial}{\partial r}\left(r \rho_{\mathrm{m}} u v\right)=-\frac{\partial p}{\partial z}+\frac{1}{r} \frac{\partial}{\partial z}\left(2 r \mu \frac{\partial u}{\partial z}\right)+\frac{1}{r} \frac{\partial}{\partial r}\left[r \mu\left(\frac{\partial u}{\partial r}+\frac{\partial v}{\partial z}\right)\right], \\
\frac{\partial}{\partial t}\left(\rho_{\mathrm{m}} v\right)+\frac{1}{r} \frac{\partial}{\partial z}\left(r \rho_{\mathrm{m}} u v\right)+\frac{1}{r} \frac{\partial}{\partial r}\left(r \rho_{\mathrm{m}} v^{2}\right)=-\frac{\partial p}{\partial r}+\frac{1}{r} \frac{\partial}{\partial z}\left[r \mu\left(\frac{\partial v}{\partial z}+\frac{\partial u}{\partial r}\right)\right]+\frac{1}{r} \frac{\partial}{\partial r}\left(2 r \mu \frac{\partial v}{\partial r}\right)-2 \mu \frac{v}{r^{2}}+\rho_{\mathrm{m}} \frac{u^{2}}{r}, \\
\frac{\partial}{\partial t}\left(\rho_{\mathrm{m}} w\right)+\frac{1}{r} \frac{\partial}{\partial z}\left(r \rho_{\mathrm{m}} u w\right)+\frac{1}{r} \frac{\partial}{\partial r}\left(r \rho_{\mathrm{m}} v w\right)=\frac{1}{r} \frac{\partial}{\partial z}\left(r \mu \frac{\partial w}{\partial z}\right)+\frac{1}{r^{2}} \frac{\partial}{\partial r}\left[r^{3} \mu \frac{\partial}{\partial r}\left(\frac{w}{r}\right)\right]-\rho_{\mathrm{m}} \frac{v w}{r}, \tag{S3}
\end{array}
$$

where $t$ is the time, $z$ represents the axial coordinates, $r$ represents the radial coordinates of the symmetry axis, $\rho_{\mathrm{m}}$ is the mixture density, $\mu$ represents the molecular viscosity, $u, v$ and $w$ denotes the axial, radial and swirl velocity components, respectively.

S2 Complete formulas of the Wall-Adapting Local Eddy Viscosity (WALE) model
The rate-of-strain tensor is calculated by:

$$
\begin{equation*}
S_{i j}=\frac{1}{2}\left(\frac{\partial \bar{u}_{i}}{\partial x_{j}}+\frac{\partial \bar{u}_{j}}{\partial x_{i}}\right) . \tag{S4}
\end{equation*}
$$

The eddy viscosity is modeled as:

$$
\begin{gather*}
\mu_{\mathrm{t}}=\rho_{\mathrm{m}} L_{\mathrm{s}}^{2} \frac{\left(S_{i j}^{\mathrm{d}} S_{i j}^{\mathrm{d}}\right)^{\frac{3}{2}}}{\left(\bar{S}_{i j} \bar{S}_{i j}\right)^{\frac{5}{2}}+\left(S_{i j}^{\mathrm{d}} S_{i j}^{\mathrm{d}}\right)^{\frac{5}{4}}},  \tag{S5}\\
S_{i j}^{\mathrm{d}}=\frac{1}{2}\left(\bar{g}_{i j}^{2}+\bar{g}_{j i}^{2}\right)-\frac{1}{3} \delta_{i j} \bar{g}_{k k}^{2}, \tag{S6}
\end{gather*}
$$

$$
\begin{gather*}
\bar{g}_{i j}=\frac{\partial \bar{u}_{i}}{\partial x_{j}}  \tag{S7}\\
L_{\mathrm{s}}=\min \left(\kappa d, C_{\mathrm{w}} V^{\frac{1}{3}}\right), \tag{S8}
\end{gather*}
$$

where $\bar{g}_{i j}$ denotes the velocity gradient tensor; $\delta_{i j}$ denotes the Kronecker symbol; $d$ is the length of the point from the closed wall; $V$ is the calculated grid volume; $\kappa$ is the Von Karman's constant of $0.41 ; C_{\mathrm{w}}$ is defined as the WALE constant of $0.325 ; S_{i j}^{\mathrm{d}}$ and $L_{\mathrm{s}}$ represent model variable defined by Eqs. (S6) and (S8), respectively.

## S3 Complete formulas of the Schnerr-Sauer model

The mass transport equation could be described as:

$$
\begin{equation*}
\frac{\partial\left(\rho_{\mathrm{v}} \alpha_{\mathrm{v}}\right)}{\partial t}+\frac{\partial\left(\rho_{\mathrm{v}} \alpha_{\mathrm{v}} u_{j}\right)}{\partial x_{j}}=\dot{m}^{+}-\dot{m}^{-} \tag{S9}
\end{equation*}
$$

where, the source term $\dot{m}^{+}$represents evaporation and $\dot{m}^{-}$represents condensation. In the bubble dynamics, the two terms are calculated through the Rayleigh-Plesset equation:

$$
\begin{cases}\dot{m}^{+}=\frac{\rho_{\mathrm{v}} \rho_{1}}{\rho} \alpha_{\mathrm{v}}\left(1-\alpha_{\mathrm{v}}\right) \frac{3}{R_{\mathrm{b}}} \sqrt{\frac{2}{3} \frac{\left(p_{\mathrm{v}}-p\right)}{\rho_{1}}}, & p \leq p_{\mathrm{v}}  \tag{S10}\\ \dot{m}^{-}=\frac{\rho_{\mathrm{v}} \rho_{1}}{\rho} \alpha_{\mathrm{v}}\left(1-\alpha_{\mathrm{v}}\right) \frac{3}{R_{\mathrm{b}}} \sqrt{\frac{2}{3} \frac{\left(p-p_{\mathrm{v}}\right)}{\rho_{\mathrm{v}}}}, & p \geq p_{\mathrm{v}}\end{cases}
$$

The bubble radius $R_{\mathrm{b}}$, vapor volume fraction $\alpha_{\mathrm{v}}$ and the bubble number density $N_{\mathrm{b}}$ are related, can be represented as:

$$
\begin{equation*}
R_{\mathrm{b}}=\left(\frac{\alpha_{\mathrm{v}}}{1-\alpha_{\mathrm{v}}} \cdot \frac{3}{4 \pi} \cdot \frac{1}{N_{\mathrm{b}}}\right)^{\frac{1}{3}} \tag{S11}
\end{equation*}
$$

In current work the bubble number density $N_{\mathrm{b}}$ is defined as a constant value of $10^{13}$.

## S4 Solution methods

| Solution method |  |
| :--- | :--- |
| Solver | Coupled |
| Gradient discretization method | Least Squares Cell Based |
| Pressure discretization method | PRESTO ! |
| Momentum discretization method | Bounded Central Differencing |
| Volume fraction discretization method | QUICK |
| Transient formulation | Bounded Second Order Implicit |

## S5 Complete grid refinement method and GCI calculation formulas

Firstly, a coarse grid, Grid 1, was generated, and then the grid spacing was halved to obtain the finer grids, Grid 2 and Grid 3. The numerical calculation solutions, $f_{1}, f_{2}$ and $f_{3}$, can be expressed using the generalized Richardson extrapolation equation as:

$$
\left\{\begin{array}{l}
f_{1}=f_{\text {exa }}+c h_{1}^{s}+o\left(h_{1}^{s+1}\right),  \tag{S12}\\
f_{2}=f_{\text {exa }}+c h_{2}^{s}+o\left(h_{2}^{s+1}\right), \\
f_{3}=f_{\text {exa }}+c h_{3}^{s}+o\left(h_{3}^{s+1}\right),
\end{array}\right.
$$

where, $f_{\text {exa }}$ denotes the calculation solution when the grid spacing approaches zero, serving as the exact solution, $c$ is a constant value, $h_{1}, h_{2}$ and $h_{3}$ represent grid spacings, and $s$ denotes the order of convergence and dependents on the calculation method.

The grid refinement ratio $r_{\mathrm{g}}$ is defined as:

$$
\begin{equation*}
r_{\mathrm{g}}=\frac{h_{1}}{h_{2}}=\frac{h_{2}}{h_{3}} . \tag{S13}
\end{equation*}
$$

Neglecting higher-order terms in Eq. (S12) and eliminating the $f_{\text {exa }}$ and $c$ :

$$
\begin{equation*}
\frac{f_{1}-f_{2}}{f_{2}-f_{3}}=\frac{h_{1}^{s}-h_{2}^{s}}{h_{2}^{s}-h_{3}^{s}} . \tag{S14}
\end{equation*}
$$

Taking the logarithm of both sides of Eq. (S14):

$$
\begin{equation*}
s=\frac{\ln \left(\frac{f_{1}-f_{2}}{f_{2}-f_{3}}\right)}{\ln r_{g}} \tag{S15}
\end{equation*}
$$

Set $\varepsilon_{1}=\left(f_{1}-f_{2}\right) / f_{1}, \quad \varepsilon_{2}=\left(f_{2}-f_{3}\right) / f_{2}$, the estimated fractional error could be expressed as,

$$
\left\{\begin{array}{l}
\frac{f_{1}-f_{\text {exa }}}{f_{1}}=\frac{f_{1}-f_{2}}{f_{1}\left[1-\left(\frac{1}{r_{g}}\right)^{s}\right]}=\frac{\varepsilon_{1} r_{\mathrm{g}}^{s}}{r_{\mathrm{g}}^{s}-1},  \tag{S16}\\
\frac{f_{2}-f_{\text {exa }}}{f_{2}}=\frac{f_{1}-f_{2}}{f_{2}\left(r_{\mathrm{g}}^{s}-1\right)}=\frac{\varepsilon_{1}}{r_{\mathrm{g}}^{s}-1}, \\
\frac{f_{3}-f_{\text {exa }}}{f_{3}}=\frac{f_{2}-f_{3}}{f_{3}\left(r_{\mathrm{g}}^{s}-1\right)}=\frac{\varepsilon_{2}}{r_{\mathrm{g}}^{s}-1} .
\end{array}\right.
$$

The Grid Convergence Index (GCI) is defined as:

$$
\begin{equation*}
\mathrm{GCI}=F_{\mathrm{s}} \cdot\left|\frac{f-f_{\text {exa }}}{f}\right|, \tag{S17}
\end{equation*}
$$

where $f$ represents the numerical calculation solution and $F_{\mathrm{s}}$ is a factor of safety, $F_{\mathrm{s}}=1.25$. GCI indicates an error band of the deviation of the solution from the exact value. Therefore, it also reflects the variation of the solution as the grid is further refined.

## S6 Nomenclature

| $\alpha$ | throttling angle, rad |
| :--- | :--- |
| $\alpha_{\mathrm{v}}$ | volume fraction number of vapor |
| $\mu_{\mathrm{l}}$ | dynamic viscosity of liquid, $\mathrm{kg} /(\mathrm{m} \mathrm{s})$ |
| $\mu_{\mathrm{m}}$ | mixing dynamic viscosity, $\mathrm{kg} /(\mathrm{m} \mathrm{s})$ |
| $\mu_{\mathrm{t}}$ | eddy viscosity, $\mathrm{kg} /(\mathrm{m} \mathrm{s})$ |
| $\mu_{\mathrm{v}}$ | dynamic viscosity of vapor, $\mathrm{kg} /(\mathrm{m} \mathrm{s})$ |
| $\rho_{\mathrm{l}}$ | density of liquid, $\mathrm{kg} / \mathrm{m}^{3}$ |
| $\rho_{\mathrm{m}}$ | mixing density, $\mathrm{kg} / \mathrm{m}^{3}$ |
| $\rho_{\mathrm{v}}$ | density of vapor, $\mathrm{kg} / \mathrm{m}^{3}$ |
| $\sigma$ | Cavitation number |
| $\tau_{i j}$ | the sub-grid scale stress |
| $L_{\mathrm{s}}$ | sealing length, mm |
| $L_{\mathrm{x}}$ | valve opening, mm |
| $\dot{m}^{+}$ | mass transfer source term connected to evaporation, $\mathrm{kg} /\left(\mathrm{m}^{3} \mathrm{~s}\right)$ |
| $\dot{m}^{-}$ | mass transfer source term connected to condensation, $\mathrm{kg} /\left(\mathrm{m}^{3} \mathrm{~s}\right)$ |
| $p$ | pressure, Pa |
| $p_{\text {in }}$ | inlet pressure, Pa |
| $p_{\text {out }}$ | outlet pressure, Pa |
| $p_{\mathrm{v}}$ | vapor pressure, Pa |
| $r_{\mathrm{a}}$ | inlet radius, mm |
| $R_{\mathrm{a}}$ | length-to-diameter ratio |
| $S_{i j}$ | rate-of-strain tensor for the resolved scale, $\mathrm{s}^{-1}$ |
| $t$ | time, s |
| $u$ | velocity, m$/ \mathrm{s}$ |
| $x$ | Cartesian coordinate, m |

