

Research Article

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Influence of ground effect on flow field structure and aerodynamic noise of high-speed trains

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Abstract: The simulation of the ground effect has always been a technical difficulty in wind tunnel tests of high-speed trains. In this paper, large eddy simulation and the curl acoustic integral equation were used to simulate the flow-acoustic field results of high-speed trains under four ground simulation systems (GSSs): “moving ground+rotating wheel”, “stationary ground+rotating wheel”, “moving ground+stationary wheel”, and “stationary ground+stationary wheel”. By comparing the fluid-acoustic field results of the four GSSs, the influence laws of different GSSs on the flow field structure, aero-acoustic source, and far-field radiation noise characteristics were investigated, providing guidance for the acoustic wind tunnel testing of high-speed trains. The calculation results of the aerodynamic noise of a 350 km/h high-speed train show that the moving ground and rotating wheel affect mainly the aero-acoustic performance under the train bottom. The influence of the rotating wheel on the equivalent sound source power of the whole vehicle was not more than 5%, but that of the moving ground slip was more than 15%. The average influence of the rotating wheel on the sound pressure level radiated by the whole vehicle was 0.3 dBA, while that of the moving ground was 1.8 dBA.

Key words: High-speed train; Aero-acoustics; Flow field structure; Large eddy simulation; Moving ground condition; Rotating wheel

1 Introduction


As high-speed ground transportation vehicles, high-speed trains inevitably encounter ground effect problems, that is, disturbance caused by the ground boundary layer to the structure of the train bottom flow fields. The accuracy of simulation of the ground effect phenomenon has a great effect on the accuracy of predictions of high-speed train aerodynamics and aerodynamic noise.

Prediction models are widely used at present to understand the nature of railway ground waves and their associated characteristics (Kouroussis et al., 2021).

Using boundary layer correction theory to correct the simulation results of aerodynamic force and aerodynamic noise of 3D complex structures such as high-speed trains is difficult. The main techniques available to eliminate the negative effects of ground effect in wind tunnel tests include the moving ground method, tangent blow-suction method, and suction method. The moving ground method is the most satisfactory for its capacity to simulate the relative motion between the train, ground, and air. Using the moving ground method, Tyll et al. (1996) estimated the aerodynamic drag coefficients of maglev to be 0.24 with moving rails and 0.17 without rails. Moreover, they pointed out that a moving track is very important to the accuracy of aerodynamic testing of a maglev train. Yi et al. (1997) investigated the ground effect of a high-speed train using the suction method. They showed that the elimination of the boundary layer can significantly increase the aerodynamic drag force coefficient, the increment of which depends on the shape of the train bottom. By comparing

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the test results from the tangent blow-suction method and the moving ground method, Kwon et al. (2001) concluded that the tangential blowing method was an effective substitute for the moving ground method for testing aerodynamics.

Xia et al. (2016) compared train wind simulation results between the stationary ground and the moving ground conditions and concluded that the stationary ground result was larger than that of the moving ground. In the following year, they constructed seven ground simulation systems (GSSs) and pointed out that the elimination of the boundary layer can significantly increase the aerodynamic drag coefficient, and that elevation of the model cannot effectively eliminate the adverse effects of the ground effect (Xia et al., 2017). Zhang et al. (2016) numerically simulated the train aerodynamic simulation results of three GSSs: “stationary ground+stationary wheel”, “moving ground+stationary wheel”, and “moving ground+rotating wheel”, and concluded that the moving ground could significantly increase the aerodynamic drag coefficient, while the rotating wheel could not. Paz et al. (2017) evaluated the influence of sleepers on the aerodynamic drag coefficient on the ground effect through the dynamic layering method. They found that the existence of sleepers had a great impact on the aerodynamic performance of the whole train, increasing the aerodynamic drag coefficient by about 15%, but had little impact on the lift coefficient. Zhu et al. (2017) numerically calculated the aerodynamic noise radiated by a wheelset of two GSS conditions: moving ground and no ground. They found that the result of moving ground was about 7 dB higher than that of no ground. The simulation results of aerodynamic noise of high-speed trains by Liu et al. (2013) showed that the result of moving ground was 4–6 dBA larger than that of stationary ground. Wen et al. (2019) found that the existence of bogies on the bottom of the train, especially the last bogie, not only enhanced the wake flow but also introduced large perturbances into the wake flow. The bogie effects were revealed through a systematic comparison of flow structures, slipstream characteristics, and aerodynamic forces between two generic train configurations (Shibo et al., 2018).

To sum up, we now have a more comprehensive and accurate understanding of how the ground effect influences the prediction accuracy of aerodynamic force, but how the ground effect influences the prediction

accuracy of aerodynamic noise is still unclear. The reason is that the current methods of eliminating the ground boundary layer cannot account for either changes in the fluctuating flow field at the train bottom or the generation of new noise. The ground is assumed to be stationary in current acoustic wind tunnel tests, and it is impossible to explore the influence laws of different GSSs on aero-acoustic tests. On the other hand, the current geometric models for numerical simulation are simplified. In this study, we constructed numerical models of four GSSs to explore the influence laws of different GSSs on the flow field structure and aerodynamic noise, and to provide guidance for the acoustic wind tunnel testing of high-speed trains.

2 Numerical simulation model

The geometric model was a 1/8-scale high-speed train model with three carriages, bogies and without a pantograph (Tan et al., 2018). The model was 7960 mm long, 4080 mm high, and 3360 mm wide. The geometric model was divided into 32 parts, i.e., Components #1–#32. The naming rules of the first 20 parts are shown in Fig. S1 of the electronic supplementary materials (ESM), and the last 12 parts correspond to the 12 wheelsets. From upstream to downstream, the wheelset was named sequentially from Component #21. For example, Component #21 corresponds to the front wheelset of Bogie 1, and Component #22 corresponds to the back wheelset of Bogie 1.

The simulated incoming velocity was 350 km/h. The surface of the train model, excluding the 12 wheelsets, was set as a non-slip boundary condition with friction. All wheelsets were set to a speed of 1817.2 or 0 rad/s. The wheel speed is derived from Section S1 of the ESM. The ground moving velocity was 350 or 0 km/h. These could be combined freely to form four GSSs: “moving ground+rotating wheel”, “stationary ground+rotating wheel”, “moving ground+stationary wheel”, and “stationary ground+stationary wheel”. For the convenience of discussion, they were named Case 1, Case 2, Case 3, and Case 4, respectively.

We used the Ansys Fluent software from the Wuxi Computing Centre (China) to carry out numerical simulations of the aerodynamic noise of the high-speed train. The unsteady calculation time step was 5×10^{-5} s, and iterated 30 steps in each time step. A total of

10000 time steps were calculated. The computation of the first 2500 time steps was to ensure full development of the turbulence flow field, and that of the remaining 7500 time steps was to extract the noise source information. Sound source data in each time step were stored, with a time span of 0.375 s and a frequency resolution of 2.7 Hz. The number of body grids was about 110 million.

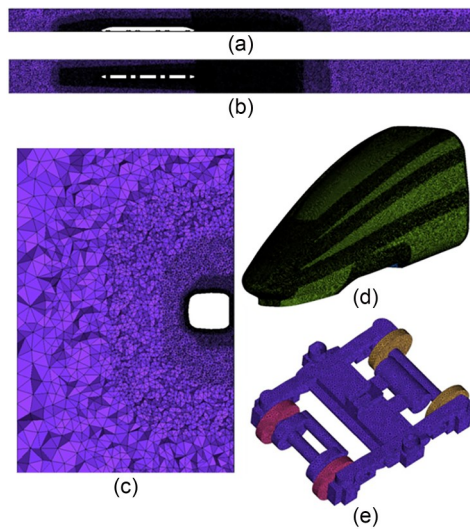


Fig. 1 Grid distribution: (a) longitudinal symmetry plane; (b) on the same elevation as the tip of the nose; (c) a cross section of the tail-flow type shoulder; (d) streamlined area; (e) bogie

3 Difference analysis of flow field structure characteristics

Through the three physical quantities of velocity amplitude, vorticity amplitude, and Q value, the differences among the four GSSs in terms of the flow field structure around the train were revealed. For convenience of description, Fig. 2 shows all sections in this study. Section a is a y -equivalent plane running

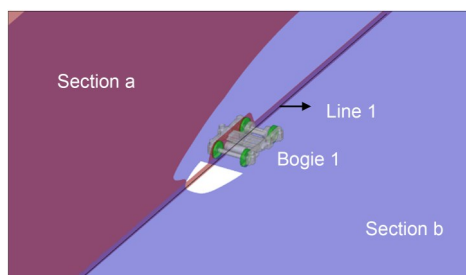


Fig. 2 Sketch map of the section

through the calculation domain and the wheel pair, with a y value of 93 mm. Section b is the z iso-surface penetrating the calculation domain, with a z value of 50 mm. Line 1 runs along the x direction and penetrates the calculation region, with a y value of 93 mm and a z value of 12.5 mm.

3.1 Flow velocity magnitude

From Figs. 3 and 4, the influence of the four GSSs on the spatial distribution of the velocity amplitude around the high-speed train can be obtained. The four cases had no influence on the velocity amplitude distribution in Region I, the upper part of the train. In Region II, near the nose of the head car, whether the wheel was rotating or not had little effect on the form of the stagnation zone, and the volume of the stagnation zone caused by “stationary ground” was larger than that caused by “moving ground”. In Region III, the area from the discharger of the head car to Bogie 1, the length of the acceleration zone under the “moving

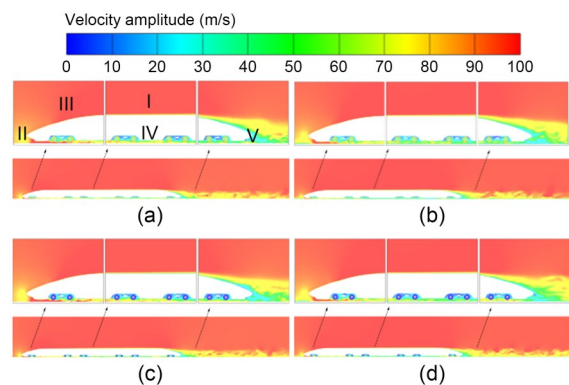


Fig. 3 Instantaneous velocity amplitude contour of Section a: (a) Case 1; (b) Case 2; (c) Case 3; (d) Case 4

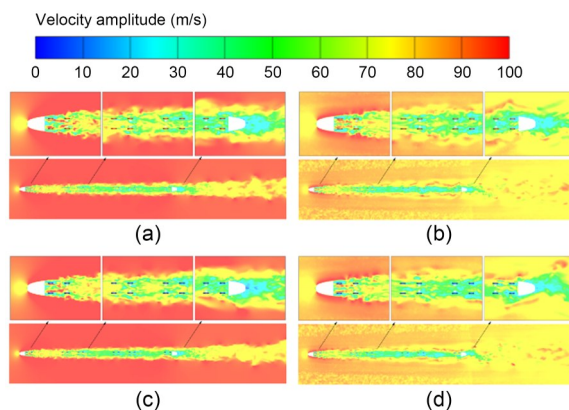


Fig. 4 Instantaneous velocity amplitude contour of Section b: (a) Case 1; (b) Case 2; (c) Case 3; (d) Case 4

ground” condition was larger than that under the “stationary ground” condition. When the wheel was rotating, the velocity amplitude inside the cavity was increased. In Region IV, the area from beneath Bogie 2 to Bogie 3, the velocity amplitude under “moving ground” was larger than that under “stationary ground”, and so was the velocity amplitude inside the cavity. In Region V, the wake area near the discharger of the tail car, the low-speed zone in “moving ground”, was smaller than that in “stationary ground”. Nevertheless, whether the wheel was rotating or not hardly affected the velocity amplitude distribution in this region. Note that the velocity amplitude distribution of the ground area under “stationary ground” was closely related to the grid scale and ground boundary layer effect.

To quantitatively analyze the velocity amplitude distribution formed in the four GSSs, Table 1 shows the mean value of the velocity amplitude of Line 1.

Table 1 Mean value of the velocity amplitude

Case No.	Mean value of velocity amplitude (m/s)
1	85.0
2	69.7
3	84.9
4	69.0

Whether the wheel was rotating or not hardly affected the velocity amplitude of Line 1, while the velocity amplitude of Line 1 on moving ground was 23% larger than that on stationary ground (Table 1).

3.2 Vorticity magnitude

From Figs. 5 and 6, we can obtain the effects of the four GSSs on the spatial distribution of vorticity amplitude around high-speed trains. The four cases had no effect on the vorticity amplitude distribution in Region I, the upper part of the train. In Region II, near the nose of the head car, the state of the wheel had little effect on the vorticity distribution. However, the vorticity distribution on “stationary ground” was uneven, but relatively uniform on “moving ground”. In Region III, the area from the discharger of the head car to Bogie 1, the “moving ground” decreased the vorticity amplitude between the train and the ballast compared with the “stationary ground”. In Region IV, the area from beneath Bogie 2 to Bogie 3, whether the ground was moving or not had little effect on the vorticity distribution in that region. In Region V, the wake area near the discharger

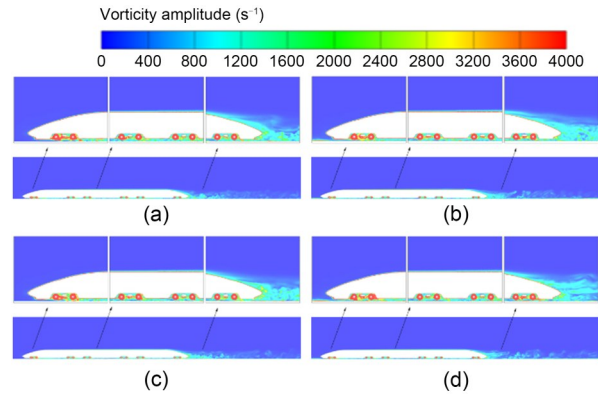


Fig. 5 Instantaneous contour plot of vorticity amplitude at Section a: (a) Case 1; (b) Case 2; (c) Case 3; (d) Case 4

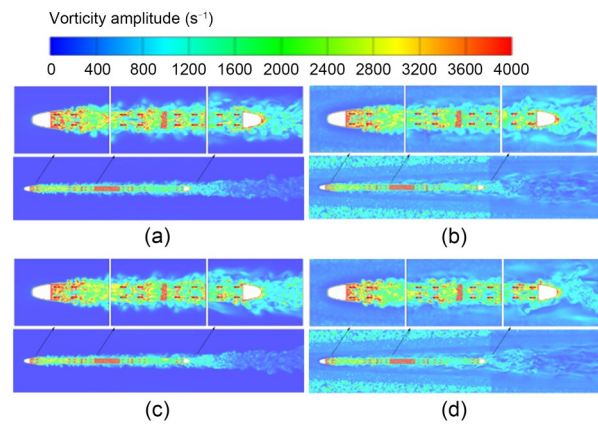


Fig. 6 Instantaneous contour plot of vorticity amplitude at Section b: (a) Case 1; (b) Case 2; (c) Case 3; (d) Case 4

of the tail car, both “stationary ground” and “moving ground” conditions could form strong vorticity separation flow at the shoulder of the tail car streamline and form a strong vortex mixing flow in the tail car cow-catcher bottom and downstream area. But the separation flow formed in the “stationary ground” condition was closer to the streamlined car body downstream, and the mixed flow formed by “stationary ground” showed large-scale rotating flow intensity and strong swing strength on both sides. Besides, whether the wheel was rotating or not hardly affected the vorticity distribution of this region. Note that the amplitude distribution of vorticity was closely related to the grid scale and the boundary layer effect. The larger the grid scale, i.e., the further away from the high-speed train area, the more uneven the cloud map of vorticity amplitude distribution, and the larger the voracity amplitude. Nonetheless, these phenomena did not appear under the “moving ground” condition.

The vorticity amplitude distribution formed in the four GSSs was analyzed. Table 2 shows the mean value of the vorticity amplitude of Line 1. Table 2 shows that whether the wheel was rotating or not hardly affected the vorticity amplitude of Line 1, while the vorticity amplitude of Line 1 on stationary ground was up to 631.3% larger than that on moving ground.

Table 2 Mean value of the vorticity amplitude

Case No.	Mean value of vorticity amplitude (s^{-1})
1	330.5
2	2262.6
3	309.4
4	2255.6

3.3 Q value

As shown in Figs. 7 and 8, we can observe the influence of the four GSSs on the spatial distribution of the scale/strength of the vortex structure around the high-speed train. The four cases had no influence on the spatial distribution of the scale/strength of vortex structure in Region I, the upper part of the train. In Region II, near the nose of the head car, the four GSSs hardly affected the vortex structure scale/strength spatial distribution. In Region III, the area from the discharger of the head car to Bogie 1, the vortex structure formed in the “moving ground” condition was larger than that in the “stationary ground” condition, and the moving speed of the vortex structure was faster along both sides. Compared with the “stationary wheel” case, the “rotating wheel” case slightly increased the scale and strength of the vortex structure inside the cavity, but decreased the number of the vortexes inside the cavity and propagated a larger group of vortexes downstream. In Region IV, the area from beneath Bogie 2 to Bogie 3, the scale of the vortex structure under “moving ground” was larger than that under “stationary ground”. There was also a stronger tendency of the vortex to move to both sides. In Region V, the wake area near the cowcatcher of the tail car, the vortex group structure rotated around the center in the “stationary ground” condition and formed a larger vortex structure on the upper part of the tail streamline than in the “moving ground” condition. However, whether the wheel was rotating or not hardly affected the formation/revolution of the vortex structure in this region. Note that the “moving ground” conveyed the vortex group structure formed in the wake region farther downstream than did the “stationary ground” condition.

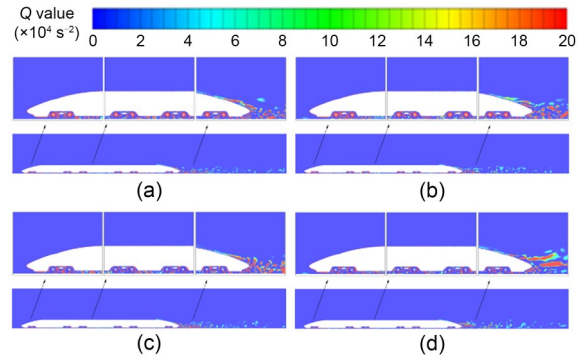


Fig. 7 Instantaneous contour plot of Q value at Section a: (a) Case 1; (b) Case 2; (c) Case 3; (d) Case 4

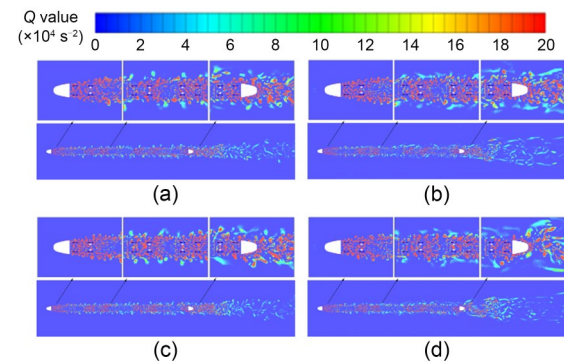


Fig. 8 Instantaneous contour plot of Q value at Section b: (a) Case 1; (b) Case 2; (c) Case 3; (d) Case 4

In summary, compared with “stationary ground”, the “moving ground” cases increased the velocity amplitude of the flow between the train and the ballast, the scale of the vortex structure, and the vortex groups transporting downstream. Furthermore, it enhanced the tendency of the vortex group under the train to move to the sides. The “rotating wheel” increased the velocity amplitude inside the cavity, the vortex amplitude and structure scale/strength, and the scale of the vortex group conveyed downstream. Also, it enhanced the shear flow release strength of the bogie guides and decreased the amount of vortex groups inside the cavity. Note that the distributions of velocity and vorticity amplitude near the ground in the “stationary ground” case were linked to the grid scale and ground boundary layer effect. However, those phenomena did not appear in the “moving ground” condition. In addition, compared with the “moving ground” condition, the vortex structure rotating around the center in the wake mixing zone under the “stationary ground” condition was more identifiable, and a large scale of vortex structure was formed above the tail streamline.

The reasons for the difference in the simulated ground flow field between the four GSSs are as follows. A boundary layer forms in the “stationary ground” and becomes thicker along the train, while the velocity inside the boundary layer is lower than the main flow, thereby preventing more air from flowing into the head car bottom and out of the tail car. Therefore, air accumulates at the head car upstream and the tail car downstream. The velocity gradient of the boundary layer is large resulting in the large vortex amplitude. But the vortex amplitude of the “moving ground” is small and conveys a strong vortex flow generated by the bogie cabin to the downstream. Therefore, the velocity amplitude of the upstream train bottom space under the “moving ground” is small, but the velocity amplitude of the downstream train bottom space is almost the same as that of the “stationary ground”.

Among the four GSSs, the “moving ground+rotating wheel” had the closest emulation of the true flow field, followed by the “moving ground+stationary wheel”, while the other two GSSs hardly met the requirements for refined simulation of high-speed train aerodynamic noise.

4 Aerodynamic noise source

The rate of change of the fluctuating pressure p' on a high-speed train can characterize the intensity of its aerodynamic noise source (Liang et al., 2020). We used p' on the high-speed train to quantitatively study the influence law of the aerodynamic noise source intensity and frequency characteristics of the high-speed train. The data for the surface-fluctuating pressure were processed in five steps. In the first step, we extracted the time-domain signal of the fluctuating pressure p on the high-speed train. In the second step, we performed the Fourier transform and conversion to 1/3 octave bands on p and obtained P . In the third step, we obtained the fluctuating pressure change rate P' according to the Fourier transform property of the derivative, i.e., Eq. (1). In the fourth step, we calculated the rate of change of the fluctuating force F' via Eq. (2). In the fifth step, we calculated the equivalent acoustic power of the sound source W_{source} .

$$P' = 2\pi f P, \quad (1)$$

$$F' = \int_s p' ds, \quad (2)$$

where p is the fluctuating pressure on the high-speed train, Pa; P is the physical quantity of the Fourier transform for p , Pa; f is the frequency, Hz; F is the fluctuating force, N; s is the area of the noise source, m²; the symbol $'$ represents the time derivative.

4.1 Intensity characteristics

The root mean square of the high-speed train surface fluctuating pressure change rate was defined as a function by

$$p'_{\text{rms}} = \left[\frac{1}{T} \int_0^T (P')^2 dt \right]^{\frac{1}{2}}, \quad (3)$$

where T is the total calculation time.

Eq. (3) can be used to characterize the average effect of sound source intensity during the sampling time. Fig. 9 shows the dipole noise source intensity distribution cloud map of the train in each of the four GSSs.

Fig. 9 shows that the strength distribution law of the vehicle dipole noise source was consistent in the four GSSs. For instance, the areas with strong dipole sources included the lower sides of each bogie cabin and the bottom of the headstock and the bogies. The strength of the dipole sources on Bogie 1 was larger than that on the other bogies. However, there was little difference in the magnitude of the train body dipole noise source intensity calculated by the four GSSs. To the train bottom downstream of the head car cow-catcher, whether the wheel was rotating or not hardly influenced the dipole noise source distribution in that region, and the area of strong dipole noise source distribution under the “moving ground” was significantly larger than that under the “stationary ground”. With regard to the six-bogie region, the strong dipole noise source distribution area was larger when the wheel was rotating. For the five-bogie region upstream, the area of strong dipole noise source distribution under the “moving ground” was significantly larger than that under the “stationary ground”. However, at the six-bogie domains, it was just the opposite. To quantitatively analyze the dipole noise sources calculated by the four GSSs, since the flow field restored by Case 1 was the most realistic, we counted the sound power of the other three GSSs based on the equivalent sound power of each component or the whole train calculated by Case 1. The component number is in Fig. S1 of the ESM.

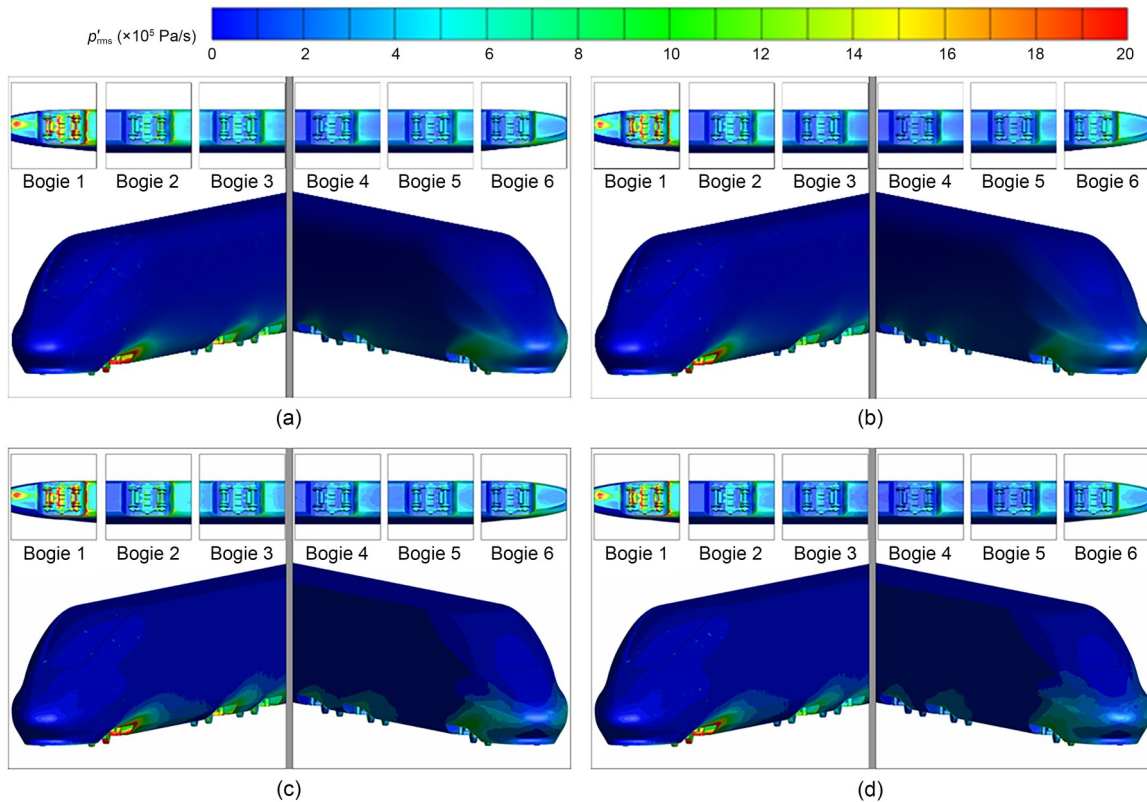


Fig. 9 Distribution cloud maps of p'_{rms} on the surface of the train: (a) Case 1; (b) Case 2; (c) Case 3; (d) Case 4

4.2 Frequency characteristics

The effects of the four GSSs on aerodynamic noise sources were discussed above in terms of aerodynamic noise source strength. We focused only on the details of the effects of the four cases on the whole car, the bottom of the head car streamlined, the bottom of the tail car streamlined, Bogie 1, Bogie 6, and the upper part of the tail car streamlined from the sound source frequency characteristics.

For the components discussed above and the whole car, we calculated the relative percentages of the other three GSSs of the sound power equivalent to Case 1 within the corresponding frequency range based on the equivalent sound power within every frequency band (1/3 octave) (Fig. 10).

With regard to the equivalent sound power of the whole car, the result calculated by Case 3 at all frequency bands was almost the same as that of Case 1, showing a maximum difference value of 1.0 dB with the corresponding frequency band center of 40 Hz. At the frequency range of [125, 10000] Hz, the results of Case 2 and Case 4 were slightly smaller than that of Case 1, with a maximum difference of 1.9 dB with the

corresponding frequency band center of 800 Hz. In the range of [20, 100] Hz, the calculation results of Case 2 and Case 4 were bigger than that of Case 1. The biggest difference was 3.2 dB and the corresponding frequency band center frequency was 50 Hz. Fig. 11 shows the percentage of the equivalent sound power result calculated by Case 2, Case 3, and Case 4 at 50 Hz relative to Case 1. The equivalent sound power at 50 Hz calculated by Case 2 and Case 4 apparently differed from that of Case 1 in the upper part of the middle car, the upper part of windshield 2, the upper part of the tail carriage and the tail car streamlined, the bottom of windshield 2, the bottom of the tail carriage and the tail car streamlined, and Bogies 4 to 6 and their wheel pairs at the downstream of the head car.

As shown in Fig. 10, the equivalent sound power of the head car streamlined bottom calculated by Case 3 was almost the same as that of Case 1 in most of the frequency bands, especially in the high-frequency range, but differed greatly in several low-frequency bands. For example, in the frequency band with a center frequency of 20 Hz the difference was 3.5 dB. The results of Case 2 and Case 4 were smaller than that of Case 1 in

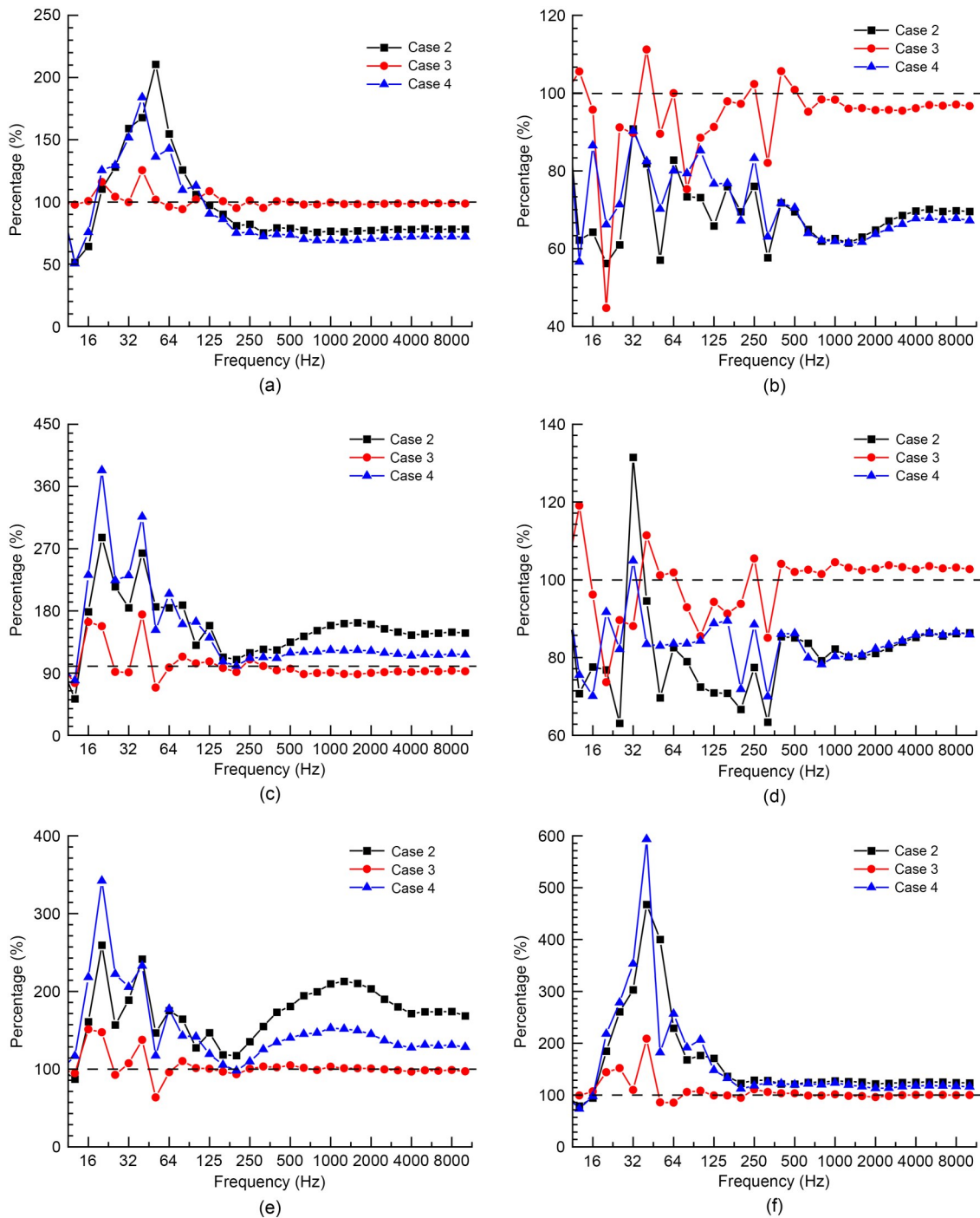


Fig. 10 Percentage of relative Case 1 calculations for each frequency band W_{source} calculated for Case 2, Case 3, and Case 4: (a) whole car; (b) head car streamlined bottom; (c) tail car streamlined bottom; (d) Bogie 1; (e) Bogie 6; (f) upper part of the tail car streamlined. The dashed line is the baseline characterized by the Case 1 calculation

all frequency bands and the maximum difference was about 2.4 dB with the frequency band center frequency being 12.5 Hz.

The equivalent sound power of the tail car streamlined bottom calculated by Case 3 was almost the same

as that of Case 1 in most frequency bands, but differed greatly in several low-frequency bands. For example, when the frequency band center frequency was 40 Hz the difference was up to 2.4 dB. In the frequency range where the center frequency was 12.5 Hz, the results of

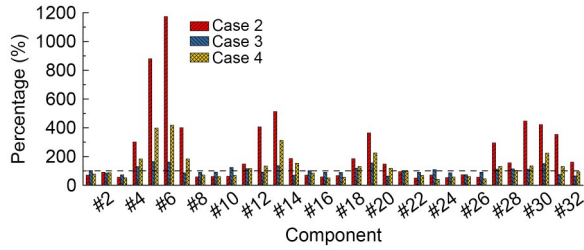


Fig. 11 Percentage of 50 Hz equivalent sound source sound calculated by Case 2, Case 3, and Case 4 relative to the Case 1 calculation result. The dashed line is the baseline characterized by the Case 1 calculation

Case 2 and Case 4 were about 2.8 and 1.0 dB smaller than that of Case 1, respectively. At other frequency ranges, the results of Case 2 and Case 4 were bigger than that of Case 1, with a maximum of about 5.8 dB and a corresponding center frequency of 20 Hz.

The equivalent sound power of Bogie 6 calculated by Case 3 was almost the same as that of Case 1 in most frequency bands, but the results were variable in several low-frequency bands. For instance, the difference was 1.8 dB when the center frequency was 16 Hz. In most frequency bands, the results of Case 2 and Case 4 were apparently larger than that of Case 1. The largest difference was about 5.3 dB with the corresponding center frequency being 20 Hz. The result of Case 2 was slightly smaller than that of Case 1 only when the frequency band center frequency was 12.5 Hz.

In general, whether the wheel was rotating and whether the ground was moving affected the equivalent sound power mainly in the frequency range below 100 Hz. Whether the ground was moving had a more significant influence on the equivalent sound power than whether the wheel was rotating. In most frequency bands, the equivalent sound power of the head car downstream calculated by the “stationary ground” was larger than that of the “moving ground”, while the components in the head car calculated by the “stationary ground” were smaller than those of the “moving ground”. We adopted the same method to analyze the percentage of the equivalent sound source, and the results are shown in Fig. S2.

5 Far field radiation noise

Combined with the mirror imaging principle, the Curl acoustic integral equation can be used to calculate the far field radiation noise considering the ground

effect. The calculation process is given in Section S2 of the ESM.

To study the characteristics of high-speed train radiated noise, we set 16 measuring points every 5 m along the train. These points were 25 m from the train’s central axis and 3.5 m above the ground. The first point was located in the 2.5 m-position downstream of the nose of the head car. The last point was located in the 2.5 m-position upstream of the nose of the tail car. The locations of these points are shown in Fig. 12. Considering the ground reflection, we set another 16 points which were symmetric about the ground according to the specular reflection principle.

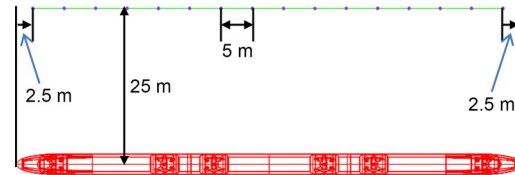


Fig. 12 Locations of the measuring points

5.1 Intensity characteristics

To study the aerodynamic noise sound pressure level (SPL), we compared the measuring point test results of the four GSSs (Fig. 13a). The difference distribution curve of whether the wheel was rotating or not is shown in Fig. 13b, and the curve of “Case 3–Case 1” represents the difference distribution curve of whether the wheel was rotating or not in the “moving ground” condition. Fig. 13c gives the difference distribution curve of whether the ground was moving or not, in which the curve marked with “Case 1–Case 2” represents the difference distribution curve in the “rotating wheel” condition and the curve “Case 3–Case 4” represents the difference distribution curve in the “stationary wheel” condition.

Fig. 13 shows that the distribution of aerodynamic noise sound pressure levels calculated by the four GSSs was consistent. But the aerodynamic noise pressure level results of the four GSSs apparently differed in their simulation values. Compared with the “stationary wheel”, the result of the “rotating wheel” was smaller in the upstream, while the result tested by the downstream points was the opposite. The mean difference value of whether the wheel was rotating or not at 16 measuring points was 0.3 dBA. Compared with “stationary ground”, the result for the “moving ground” of each point was larger, especially the points in the

upstream, while in the downstream the difference was smaller. The average difference between whether the ground was moving or not at the 16 points was 1.8 dBA. Whether the ground was moving or not had a more significant influence on the aerodynamic noise sound pressure level radiated by the whole train.

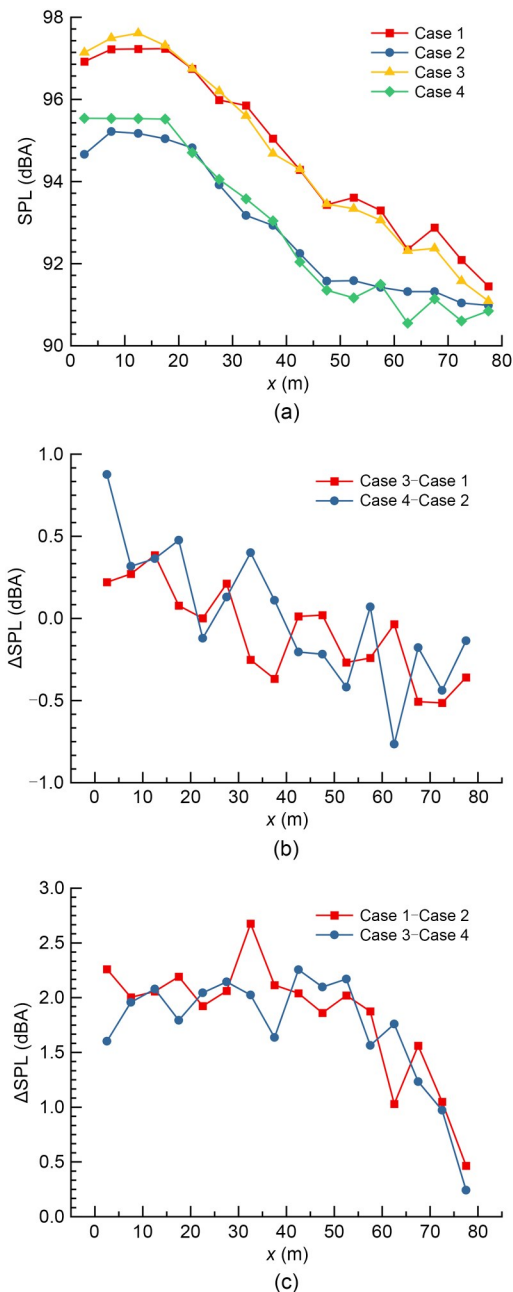


Fig. 13 Aerodynamic noise sound pressure level radiated by the vehicle to each measuring point: (a) results of the four GSSs; (b) difference between whether the wheel pair rotated or not; (c) difference between whether the ground was moving or not. SPL represents the A-weighted sound pressure level (dBA) and the reference quantity is 2×10^{-5} Pa

To understand the influence of the four GSSs on the radiated aerodynamic noise pressure level of each part, we first calculated the A-weight aerodynamic noise SPL radiated by each component at these measuring points. The calculation process is given in Section S3 of the ESM. Fig. 14 shows the difference of $SPL_{average}$ between the “rotating wheel” and “stationary wheel” for each part. Fig. 15 shows the difference of $SPL_{average}$ between the “moving ground” and “stationary ground” for each component.

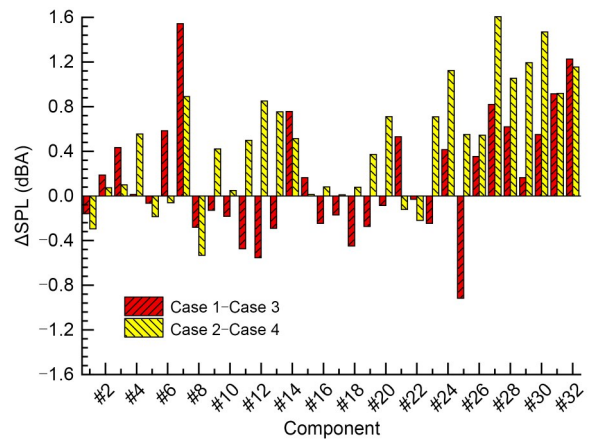


Fig. 14 Histogram of the difference in $SPL_{average}$ between the “rotating wheel” and “stationary wheel”

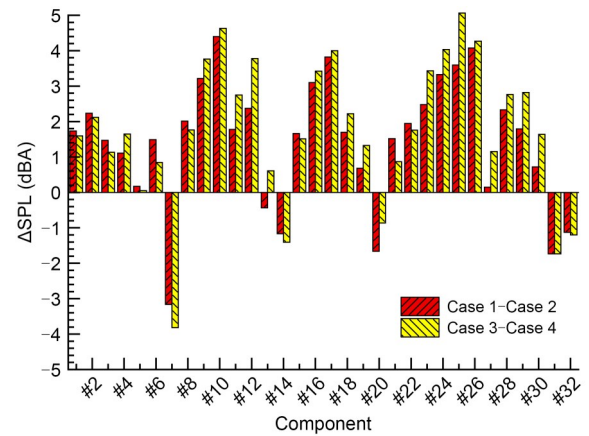


Fig. 15 Histogram of the difference in $SPL_{average}$ between the “moving ground” and “stationary ground”

Fig. 14 shows that the effect of whether the wheel was rotating or not on the $SPL_{average}$ radiated by most components was slight. For example, in the head car streamlined part, the difference in $SPL_{average}$ between the “rotating wheel” and “stationary wheel” conditions was -0.2 dBA. But it had a significant effect on the downstream wheel pairs and the upper/bottom of the tail

car streamlined. For instance, the difference in $SPL_{average}$ between the “rotating wheel” and “stationary wheel” conditions was 1.5 dBA in the upper part of the tail car streamlined and 0.9 dBA in the rear wheel pair of Bogie 6. Note that the component which was slightly affected by whether the wheel was rotating or not was concentrated upstream of the tail car, especially those parts on the train, and the velocity amplitudes of the flow fields around them were large. In contrast, the velocity amplitudes of the flow fields around those that were significantly affected by the condition of the wheel were small. Otherwise, whether the wheel rotated or not had a more significant effect on the $SPL_{average}$ radiated by the downstream wheel pairs in the “stationary ground” than in the “moving ground” condition.

As Fig. 15 shows, the result $SPL_{average}$ radiated by “moving ground” was apparently larger than that for “stationary ground” for most of the components. Note that the effect of whether the ground was moving or not on the $SPL_{average}$ radiated by the components in the train bottom was stronger than the part on the train. Compared with the “stationary ground”, the result of $SPL_{average}$ calculated by the “moving ground” was larger in the tail car streamlined upstream, but smaller in the tail car streamlined. The effect of whether the ground was moving or not on the $SPL_{average}$ radiated by each component was stronger than that of the wheel motion condition.

5.2 Frequency characteristics

In this section, we discuss the influence of the four GSSs on the characteristics of the aerodynamic noise radiated by the whole car and its components. We discuss the influence of the aerodynamic noise spectrum characteristics in detail for the whole car, the bottom of the head/tail car streamlined.

Fig. 16 shows that the $SPL_{average}$ spectral curves of the four GSSs were consistent. In contrast, the $SPL_{average}$ values calculated by the four GSSs of the whole car in each frequency band were apparently different. For instance, the difference value between the “rotating wheel” and “stationary wheel” was -3.5 dBA at 25 Hz and 5.9 dBA at 50 Hz. Compared with the “stationary ground”, the $SPL_{average}$ radiated by the whole car under the “moving ground” condition was slightly larger when the frequency was above 100 Hz, while the result of “moving ground” was smaller at most of the frequencies below 100 Hz. For example, when the frequency

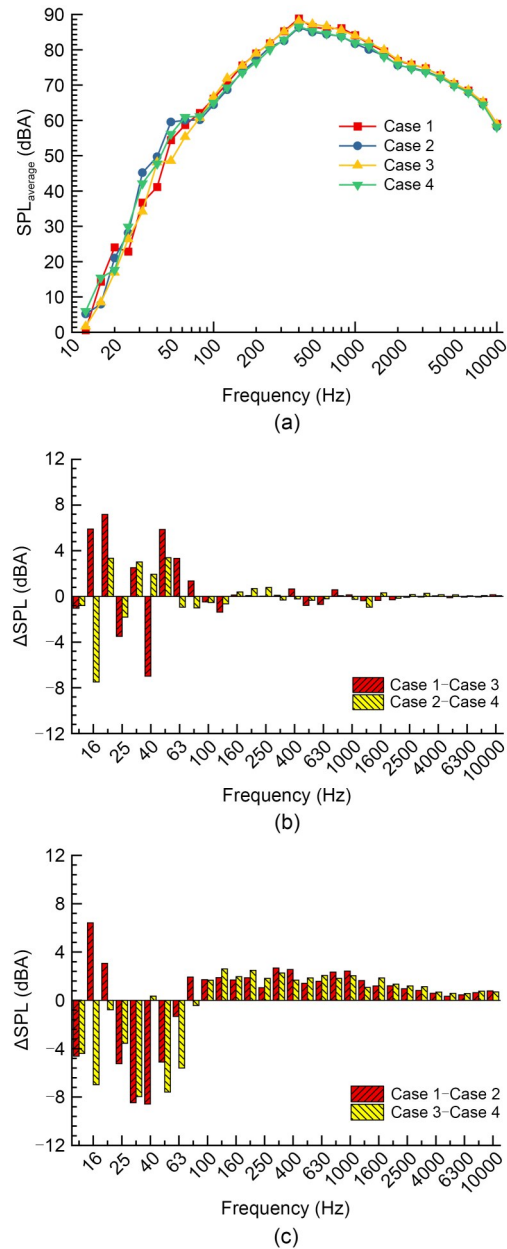


Fig. 16 $SPL_{average}$ spectrum curve radiated by the whole car: (a) result of the four GSSs; (b) value of the difference between “rotating wheel” and “stationary wheel”; (c) value of the difference between “moving ground” and “stationary ground”

was 31.5 Hz, the difference value between “moving ground” and “stationary ground” was -8.5 dBA. Compared with whether the wheel was rotating or not, the influence of whether the ground was moving or not was more apparent in the $SPL_{average}$ spectrum curve radiated by the whole car.

We adopted the same data processing method to deal with the Bogies 1 and 6 and the upper part of the

tail car streamlined as we used to deal with the $SPL_{average}$ spectral curve radiated by the whole train. The results are plotted in Figs. S3–S5.

According to Figs. 17 and 18, the spectrum curves of the above-mentioned components calculated by the four GSSs were identical, showing the distribution laws as small at both ends and large in the middle. However, the radiation of the components calculated by the four GSSs in each frequency band was quite different.

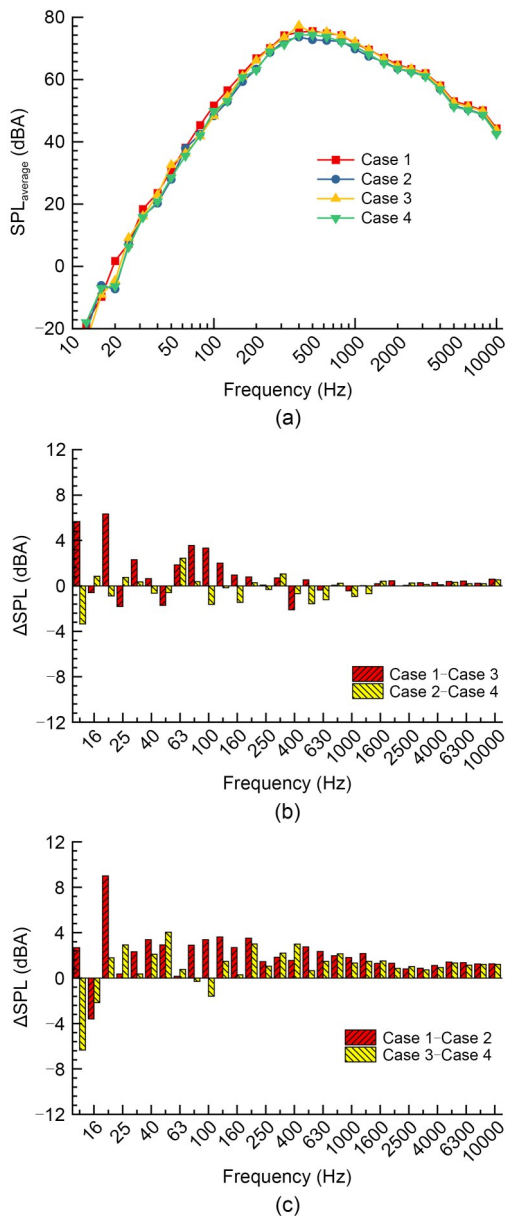


Fig. 17 $SPL_{average}$ spectral curve radiated by the bottom of the head car streamlined: (a) result of the four GSSs; (b) value of the difference between “rotating wheel” and “stationary wheel”; (c) value of the difference between “moving ground” and “stationary ground”

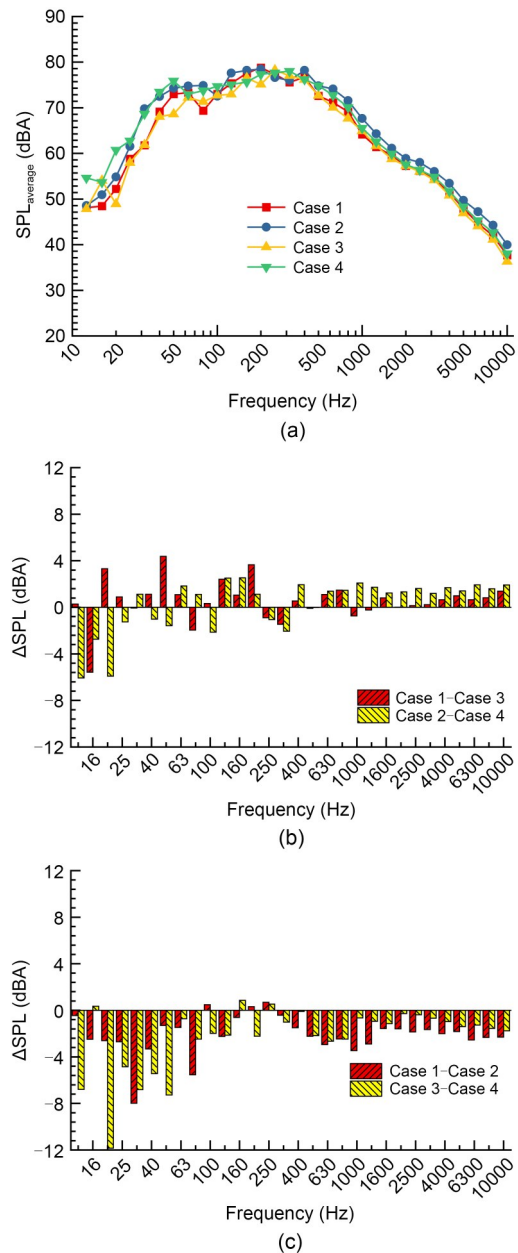


Fig. 18 $SPL_{average}$ spectral curve radiated by the bottom of the tail car streamlined: (a) result of the four GSSs; (b) value of the difference between “rotating wheel” and “stationary wheel”; (c) value of the difference between “moving ground” and “stationary ground”

The rotating wheel increased the $SPL_{average}$ radiated by the components in most frequency bands (especially the downstream components) compared with the “stationary wheel” condition. Whether the wheel was rotating or not had a bigger influence on the frequency band below 100 Hz than that above 100 Hz. The wheel condition had a more significant influence on the $SPL_{average}$ radiated by the components downstream than upstream.

In most frequency bands, the SPL_{average} radiated by the head car streamlined bottom in the “moving ground” condition was larger than that of the “stationary ground” condition, while in the bottom of the tail car streamlined were smaller in “stationary ground”. Whether the ground was moving or not had a significant influence on the SPL_{average} radiated by the components discussed above when the frequency was below 100 Hz. Whether the ground was moving or not had a more significant influence on the SPL_{average} spectrum curve radiated by the components mentioned above than whether the wheel was rotating or not.

The “moving ground” condition increased the velocity amplitude of the flow beneath the train, the scale of the vortex, the number of vortex groups transported downstream, and the tendency of the vortex group beneath the train to move to both sides of the train. This phenomenon did not appear in the “moving ground” condition, which is why the aerodynamic noise sources of the high-speed train and the noise intensity and frequency characteristics of the far field radiation were visibly different in the “stationary ground” and “moving ground” conditions.

In summary, the distribution law of the SPL_{average} spectrum curve calculated by the whole train and the components discussed above calculated by the four GSSs were consistent. However, the value of SPL_{average} radiated by the whole train and the components mentioned above varied greatly in each frequency band. The frequency band affected by whether the wheel was rotating or whether the ground was moving was concentrated below 100 Hz. Whether the ground was moving had a more significant effect on the SPL_{average} spectrum radiated by the whole train and the components mentioned above than whether the wheel was rotating. The SPL_{average} radiated by the downstream of the head car components was larger than that of the “moving ground” condition, while the SPL_{average} radiated by the components in the head car were smaller than that of the “moving ground” condition.

6 Conclusions

In this study, we investigated the influence of four GSSs on the flow structure of a high-speed train, the sources of noise, and the far field radiation noise. Whether the wheel rotated or not and whether the ground

moved or not hardly affected the flow field above the roof of the train. They affected mainly the flow field at the bottom of the train, where they had little effect on the aero-acoustic performance of the upper parts of the train, but had a significant effect on the aero-acoustic performance of the lower parts of the train. The effect of the wheels on the equivalent sound power of the whole train was less than 5%, but the effect of the ground movement condition on the equivalent sound power of the train was more than 15%. The average level of influence of rotation or non-rotation of the wheelset on the sound pressure level of aerodynamic noise was 0.3 dBA, while the level of influence of ground movement or not was 1.8 dBA. Wheelset rotation and ground movement mainly affected the equivalent sound power of the sound source and the sound pressure level below 100 Hz.

Considering that “wheelset rotation” increased the instability of numerical calculation, we suggest that only “ground slip” should be considered in the numerical calculation of the aerodynamic noise of high-speed trains. Considering that the ground was stationary in the acoustic wind tunnel experiment, we suggest that the experimental results of the high-speed train should be corrected appropriately in an acoustic wind tunnel.

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Author contributions

Xiaoming TAN designed the research. Xiaohong ZHANG and Zhigang YANG processed the corresponding data. Xiaoming TAN wrote the first draft of the manuscript. Linli GONG helped to organize the manuscript. Linli GONG revised and edited the final version.

Conflict of interest

Xiaoming TAN, Linli GONG, Xiaohong ZHANG, and Zhigang YANG declare that they have no conflict of interest.

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Electronic supplementary materials

Sections S1–S3, Figs. S1–S5, and Eqs. (S1)–(S3)