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 E-mail: jzus@zju.edu.cn



## Correspondence:

# Wideband circular-polarized transmitarray for generating a high-purity vortex beam\*

Liangjie QIU<sup>1,2,3,4</sup>, Xiuping LI<sup>†1,2,3,4</sup>, Zihang QI<sup>1,2,3,4</sup>, Wenyu ZHAO<sup>1,2,3,4</sup>, Yuhan HUANG<sup>5</sup>

<sup>1</sup>State Key Laboratory of Information Photonics and Optical Communications, Beijing 100876, China

<sup>2</sup>MOE Key Laboratory of Universal Wireless Communications, Beijing 100876, China

<sup>3</sup>Beijing Key Laboratory of Work Safety Intelligent Monitoring, Beijing 100876, China

<sup>4</sup>School of Electronic Engineering, Beijing University of Posts and Telecommunications, Beijing 100876, China

<sup>5</sup>Beijing Institute of Spacecraft System Engineering, Beijing 100094, China

<sup>†</sup>E-mail: xpli@bupt.edu.cn

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In this correspondence, a wideband circular-polarized (CP) transmitarray (TA) in the Ka-band is presented for generating a high-purity vortex beam. The proposed element is composed of two identical combinations separated by an air layer. The subwavelength structure and double-resonance characteristics ensure the stable phase shift of the element within the 1-dB transmission bandwidth of 28.4%. A square aperture TA fed by a horn antenna is fabricated and measured. Owing to the honeycomb arrangement of elements, the mode purity of  $l=-1$  is  $>0.93$  in a wide band from 28.5 to 38 GHz. The measured peak gain is 22.3 dBic, and the 3-dB axial ratio bandwidth is 27.6%. The performance of the proposed antenna demonstrates its potential for high-capacity wireless communication and high-quality radar imaging.

## 1 Introduction

With the development of wireless communication, the demand for wider bandwidth and higher data rates has significantly increased. The vortex beam carrying orbital angular momentum (OAM) is helically distributed, and was first proved in the field of optics (Allen et al., 1992). Theoretically, different OAM modes are orthogonal to each other, which enables multichannel wireless communication (Tamburini et al., 2012; Yan et al., 2014) and high-resolution radar imaging (Liu K et al., 2015; Liu HY et al., 2020).

Various approaches for generating vortex beams in microwave fields have been reported, including the patch antenna, uniform circular array, metasurface, and so on. The OAM patch antenna operating in high-order mode is good at obtaining high-purity characteristics. Nevertheless, the single-feed patch antenna (Li et al., 2020) faces the shortcomings of low gain and narrow bandwidth, whereas the feed network of the multiple-feed patch (Huang YH et al., 2019) is complex. A uniform circular array is a conventional method for generating vortex beams. However, the limited space between the elements dramatically increases the complexity of the feed network,

<sup>‡</sup> Corresponding author

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ORCID: Liangjie QIU, <https://orcid.org/0009-0004-2820-4221>;  
 Xiuping LI, <https://orcid.org/0000-0003-4350-9651>

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especially for high-purity vortex beams (Ma et al., 2021). In contrast to the above antennas, the OAM metasurface is fed by other source antennas and relies on adjusting the phase of each element. Thus, the metasurface has several advantages including simple feed and high degree of design freedom.

The metasurface falls into two main classifications, i.e., reflectarray (RA) and TA. In comparison to RA, the effect of feed-blockage is reduced in TA (Wu Z et al., 2005). So far, many TA elements have been proposed (Jiang et al., 2018; Ran et al., 2020; Veljovic and Skrivervik, 2020; Zhang FH et al., 2020; Zhang XL et al., 2020; Shahmirzadi et al., 2021). In Ran et al. (2020), the 1-dB transmission bandwidth was 25%, but the high profile was not suitable for wireless communication systems. In Jiang et al. (2018), Zhang FH et al. (2020), and Zhang XL et al. (2020), the elements consisted of fewer dielectric layers, but the transmission loss was  $>1$  dB within the operating band. In Veljovic and Skrivervik (2020) and Shahmirzadi et al. (2021), the element achieved a low profile and high transmission simultaneously, but the 1-dB transmission bandwidth was  $<7\%$  and limited the TA's bandwidth.

To generate the desired vortex beam, the phase shift of the element should meet the  $360^\circ$  requirement. So far, there are two main phase modulation methods commonly used in TA elements, i.e., the size-varying method (Akram MR et al., 2019) and the element rotation method (Zhang FH et al., 2019). The latter method is based on the Pancharatnam–Berry phase mechanism, which is also used to generate optical OAM beams (Wang et al., 2020; Wu YH et al., 2022). Notably, phase curve linearity is key to generating high-purity vortex beams. de Cos et al. (2011) revealed that the hexagonal element is smaller than the square element under the same operating bandwidth, and the phase linearity is also better. Huang YH et al. (2021) also verified this performance on size-varying elements. In addition, using the subwavelength element (Zhang FH et al., 2019) instead of the conventional half-wavelength or larger element (Bi et al., 2018; Lin et al., 2020; Wu GB et al., 2020) is beneficial for mode purity. In Huang HF and Li (2019), the period of elements was  $0.38\lambda_0$  ( $\lambda_0$  is the free-space wavelength at the center frequency), and the mode purity achieved 0.93 at the center frequency. Unfortunately, although the element was quarter wavelength (Akram Z et al., 2019),

the mode purity suffered a significant decrease from 0.96 at the center frequency to 0.72 at the edge of the operating band. So, maintaining high purity in a wide band is still challenging.

The objective of this correspondence is to propose a wideband TA element for generating a high-purity vortex beam. The TA element achieves a good balance between the bandwidth and profile. The subwavelength elements are arranged in a honeycomb shape to provide more accurate phase compensation, thereby obtaining higher mode purity. As proof of the theoretical analysis, a CP TA in the Ka-band is fabricated and measured.

## 2 Wideband circular-polarized element design

### 2.1 Element rotation principle

The element rotation principle is based on the Pancharatnam–Berry phase mechanism. Assume that there is a left-handed circularly polarized (LHCP) wave traveling in the  $+z$ -direction illuminating the element, and that the CP transmission and reflection coefficients are expressed as (Xu et al., 2017; Akram MR et al., 2020)

$$t_{LR} = \frac{1}{2}[(t_{xx} - t_{yy}) - j(t_{xy} + t_{yx})]e^{-j2\psi}, \quad (1)$$

$$r_{LL} = \frac{1}{2}[(r_{xx} - r_{yy}) - j(r_{xy} + r_{yx})]e^{-j2\psi}, \quad (2)$$

$$t_{LL} = \frac{1}{2}[(t_{xx} + t_{yy}) - j(t_{xy} - t_{yx})], \quad (3)$$

$$r_{LR} = \frac{1}{2}[(r_{xx} + r_{yy}) - j(r_{xy} - r_{yx})], \quad (4)$$

where “L” and “R” stand for LHCP and right-handed circularly polarized (RHCP) waves respectively,  $\psi$  is the azimuthal rotation angle of the element,  $t_{LR}$  is the transmission coefficient from the LHCP wave to the RHCP wave,  $r_{LL}$  is the reflection coefficient from the LHCP wave to the LHCP wave,  $t_{LL}$  is the transmission coefficient from the LHCP wave to the LHCP wave, and  $r_{LR}$  is the reflection coefficient from the LHCP wave to the RHCP wave. Because there is a polarization conversion in the transmission process, the co-polarization of the transmitted wave is RHCP. As can be seen from Eq. (1), when rotating the element in the same direction, the phase of  $t_{LR}$  is twice of the rotation angle  $\psi$ . The expressions for the linearly transmitted and reflected waves in relation to

the linear-polarized component of the incident wave can be described as

$$\begin{bmatrix} E_{tx} \\ E_{ty} \end{bmatrix} = \begin{bmatrix} t_{xx} & t_{yx} \\ t_{xy} & t_{yy} \end{bmatrix} \begin{bmatrix} E_{ix} \\ E_{iy} \end{bmatrix},$$

$$\begin{bmatrix} E_{rx} \\ E_{ry} \end{bmatrix} = \begin{bmatrix} r_{xx} & r_{yx} \\ r_{xy} & r_{yy} \end{bmatrix} \begin{bmatrix} E_{ix} \\ E_{iy} \end{bmatrix},$$

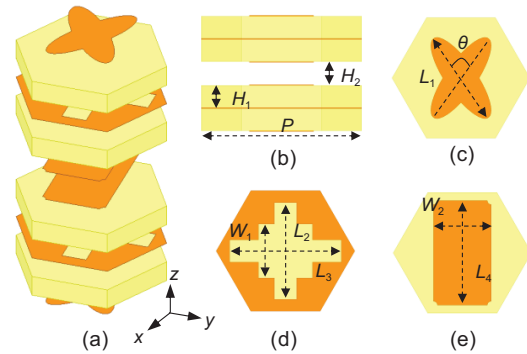
where  $x$  and  $y$  are linear polarizations,  $\mathbf{E}_t$  and  $\mathbf{E}_r$  are the vectors representing the transmitted and reflected waves respectively, and  $\mathbf{E}_i$  is the vector of the incident wave.

### 2.2 Configuration and performance

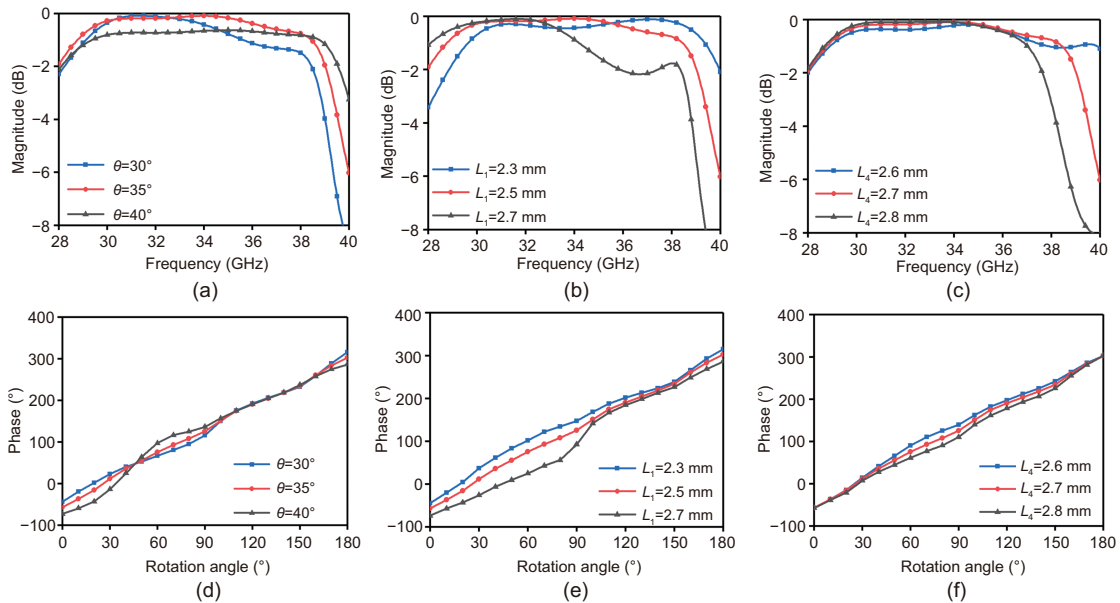
The geometry of the TA element is shown in Fig. 1. The TA element consists of two identical combinations separated by an air layer. A combined metal layer form contains an X-shaped patch, a hexagonal patch etched cross slot, and a rectangular patch. All of the dielectric layers use F<sub>4</sub>BM220 (dielectric constant  $\epsilon_r=2.2$ , loss tangent  $\tan \delta=0.001$ ) with a thickness of  $H_2=0.5$  mm. The period of the element is 3.6 mm (corresponding to  $0.40\lambda_0$ , where  $\lambda_0$  is the free-space wavelength at 33.5 GHz). To achieve a wide band and a phase shift curve with good linearity, the influence of structure parameters on the transmission coefficient is investigated, as shown in Fig. 2. Finally, the angle  $\theta$  of the X-shaped patch is  $35^\circ$ . Other detailed parameters are

set as follows (unit: mm):  $L_1=2.5$ ,  $L_2=2.5$ ,  $L_3=2.9$ ,  $L_4=2.7$ ,  $W_1=0.7$ ,  $W_2=1.5$ ,  $H_1=0.5$ . In addition, the proposed element has a relatively low profile and the total thickness of the element is only 2.5 mm (corresponding to  $0.28\lambda_0$ ).

The magnitudes of transmission and reflection coefficients are shown in Fig. 3. It can be seen that the magnitude of  $t_{LR}$  is maintained above  $-1$  dB from 29 to 38.5 GHz. Within this frequency range, the magnitude of the cross-polarized transmission coefficient is below  $-15$  dB, indicating the high



**Fig. 1** Geometry of the proposed transmitarray (TA) element: (a) exploded view; (b) front view; (c) top view of the X-shaped patch; (d) top view of the hexagonal patch; (e) top view of the rectangular patch. Yellow denotes the dielectric layers and orange denotes the metal layers. References to color refer to the online version of this figure



**Fig. 2** Magnitude and phase of co-polarized transmission coefficients versus frequency and rotation angle with different  $\theta$  (a, d),  $L_1$  (b, e), and  $L_4$  (c, f)

purity of the handedness-preserved CP transmission. Moreover, the magnitude of the reflection coefficient is below  $-10$  dB, which avoids a high back lobe level. Based on these results and the comparison from Table 1, it follows that the proposed element achieves a good balance between low profile and wide band.

The stability of the transmitted phase shift is also an important index and affects the phase compensation accuracy of TA. From Fig. 4a, a full  $360^\circ$  phase range can be covered in a wide band, and the magnitude of the co-polarized transmission coefficient is simultaneously higher than  $-1$  dB. In addition, when the rotation angle is gradually increased from  $0^\circ$  to  $180^\circ$  with an interval of  $30^\circ$ , the corresponding transmission phase responses are approximately double of the rotation angles, which agrees with the aforementioned theory. As shown in Fig. 4b, when the oblique incidence angle is  $<30^\circ$ , the phase curve is almost linear and the high transmission is maintained in the rotation process. Comparing Fig. 4a with Fig. 4c, the phase linearity of the hon-

eycomb element is significantly better than that of the square element.

### 3 Vortex beam transmitarray design

Based on the phase characteristics of a vortex beam carrying OAM (Allen et al., 1992), the required phase shift  $\Psi_{mn}$  of each element for the desired beam can be determined as

$$\Psi_{mn} = k_0 \cdot (|\mathbf{r}_{mn} - \mathbf{r}_f| - \mathbf{r}_{mn} \cdot \mathbf{u}) + l \cdot \arctan \frac{Y_{mn}}{X_{mn}}, \quad (5)$$

where  $k_0$  is the wave number in the free space corresponding to the operating frequency,  $\mathbf{r}_{mn}$  is the position vector of the element on the  $m^{\text{th}}$  row and  $n^{\text{th}}$  column,  $\mathbf{r}_f$  is the vector from the phase center of the feed to the TA's center,  $\mathbf{u}$  is the direction vector of the desired beam,  $l$  represents the OAM mode, and  $(X_{mn}, Y_{mn})$  is the coordinate of the element on the  $m^{\text{th}}$  row and  $n^{\text{th}}$  column on the beam's normal plane.

The influence of the arrangement method on OAM mode purity is analyzed. Based on the phase compensation in Eq. (5), two square aperture TAs arranged in square and honeycomb shapes are designed separately. The OAM mode purity is calculated using the numerical Fourier transform of the aperture phase function. The relationship between the OAM spectrum  $A_{l_n}$  and the sampling phase  $\Psi(\psi)$  is given by Yao et al. (2006) as follows:

$$A_{l_n} = \frac{1}{2\pi} \int_0^{2\pi} \Psi(\psi) e^{-jl_n\psi} d\psi, \quad (6)$$

$$\text{Mode purity} = \frac{|A_{l_n}|}{\sum_{m=-\infty}^{+\infty} |A_{l_m}|}. \quad (7)$$

As shown in Fig. 5 and Table 2, when the effective radiation areas are both  $10\lambda_0 \times 10\lambda_0$  and the periods

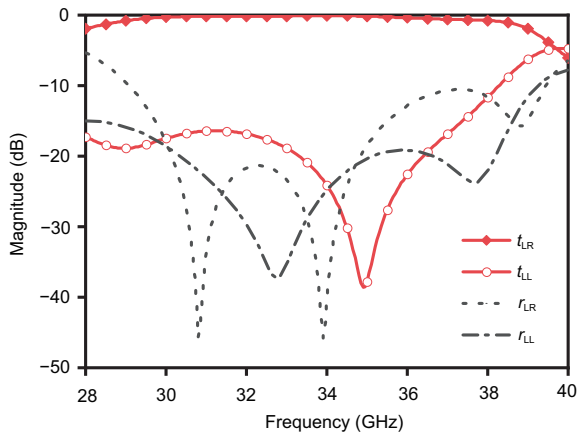


Fig. 3 Magnitude of transmission and reflection coefficients versus frequency under the normal incidence

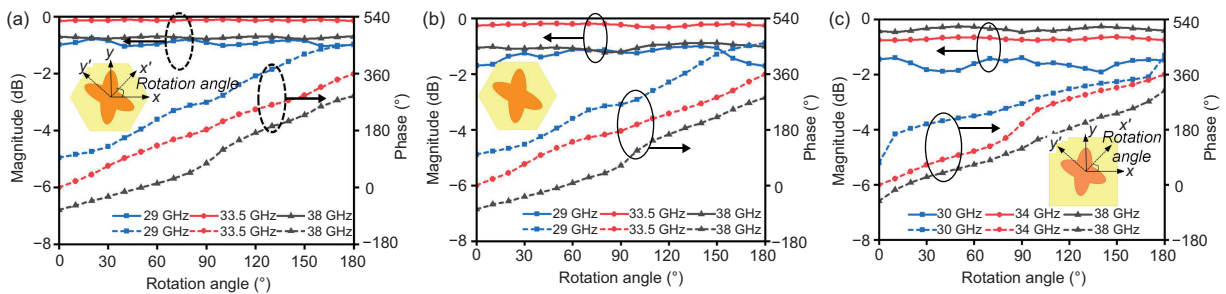


Fig. 4 Magnitude and phase of co-polarized transmission coefficient versus rotation angle of the element at different frequencies: (a) under normal incidence; (b) oblique incidence angle= $30^\circ$ ; (c) under normal incidence

**Table 1 Comparison with other reported transmitarray (TA) elements based on the element rotation principle**

Reference	$f_0$ (GHz)	Number of dielectric layers	Profile ( $\times\lambda_0$ )	Transmission bandwidth*
Shahmirzadi et al. (2021)	34	2	0.18	3.5% ( $>-1$ dB)
Veljovic and Skrivervik (2020)	24.6	2	0.13	6.0% ( $>-1$ dB)
Jiang et al. (2018)	30	2	0.31	9.1% ( $>-1.9$ dB)
Zhang XL et al. (2020)	20	2	0.21	10.0% ( $>-2.7$ dB)
Zhang FH et al. (2020)	10	3	0.14	17.5% ( $>-3$ dB)
Ran et al. (2020)	15	4	0.66	25.0% ( $>-1$ dB)
<b>This work</b>	<b>33.5</b>	<b>4</b>	<b>0.28</b>	<b>28.4%</b> ( $>-1$ dB)

\* Transmission bandwidth is the frequency range in which the magnitude of the co-polarized transmission coefficient is greater than the value in brackets

**Table 2 Comparison with other transmitarrays (TAs) based on different arrangement methods**

Arrangement method	Number of elements	Mode purity	OAM bandwidth (%)
Square	625	$>0.88$	23.5
<b>Honeycomb</b>	<b>941</b>	<b><math>&gt;0.93</math></b>	<b>28.4</b>

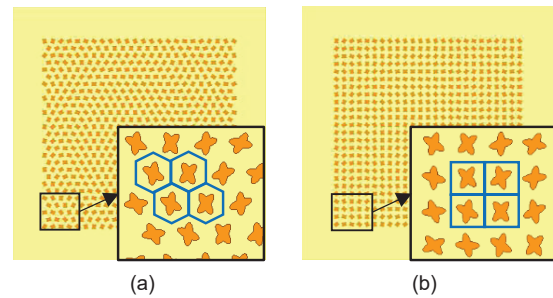
Aperture size:  $10\lambda_0 \times 10\lambda_0$ . OAM: orbital angular momentum

of the elements are both  $0.4\lambda_0$ , the honeycomb TA can accommodate more elements compared with the square TA. Moreover, the phase linearity of the honeycomb elements is better. Thus, the phase compensation accuracy is improved and higher mode purity is obtained.

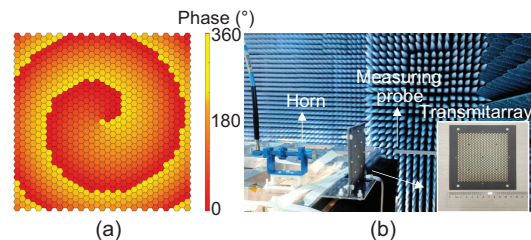
Therefore, the honeycomb arrangement method is chosen. The effective radiation area of the honeycomb TA is  $90\text{ mm} \times 90\text{ mm}$ . An LHCP horn antenna with a maximum gain of 15.2 dBic is used as a feeding source. The radiation pattern of the horn is approximately represented as a  $\cos^{q_f}(\theta)$  function (feed pattern factor  $q_f=13.5$ ), and the axial ratio is  $<3$  dB in the whole Ka-band. To achieve a good trade-off between the spillover loss and the amplitude tapering across the TA, the  $F/D$  (focal-diameter ratio) is set as 1.0. Consequently, the maximum oblique incidence angle is  $26.6^\circ$ , and the proposed TA elements can be applied. The phase distribution for  $l=-1$  is shown in Fig. 6a.

## 4 Results and discussion

A TA for generating a vortex beam is simulated, fabricated, and measured by a planar near-field measurement system (Fig. 6b). A  $90\text{ mm} \times 90\text{ mm}$  scanning plane is set 300 mm away from the TA with a measuring step of 5 mm. The simulation and measurement results in the near field are provided in Fig. 7. It is clear that the electric field (E-field) phase is a standard anticlockwise spiral with  $l=-1$ ,



**Fig. 5 Schematic of transmitarray (TA) using different arrangement methods: (a) honeycomb; (b) square**



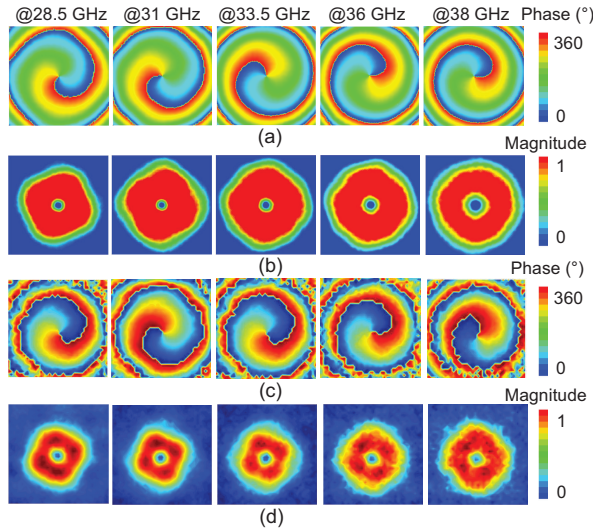
**Fig. 6 Phase distribution of transmitarray (TA) (a) and configuration of the planar near-field measurement system (b)**

and the magnitude has a central hole from 28.5 to 38 GHz. Moreover, as depicted in Fig. 8, the OAM mode purity of  $l=-1$  is  $>0.93$  in the range of 28.5–38 GHz, and is still high at the edge of the operating band, which illustrates the good quality of the generated vortex beam. In other words, the mode purity bandwidth is 28.4%.

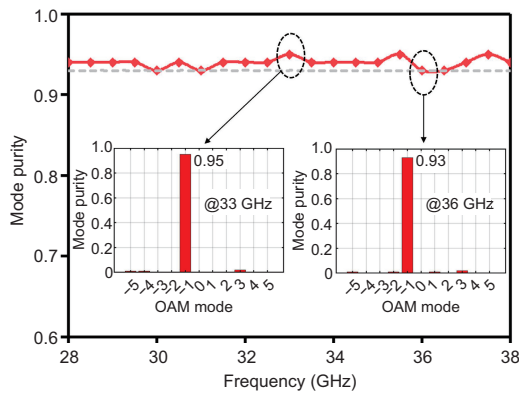
The results of gain and axial ratio are plotted in Fig. 9. The measured peak gain of the proposed antenna is 22.3 dBic at 37.5 GHz and the 1-dB gain bandwidth is 17.1%. The 3-dB axial ratio bandwidth

is 27.6% from 28 to 37 GHz. A slight discrepancy between the simulation and measurement results is mainly due to the increase in the air gap between the dielectric layers caused by the fabrication tolerance.

The structural characteristics and performance



**Fig. 7** Results of E-field phase and magnitude distributions in the near field for  $l=-1$  at different frequencies: (a) simulated phase; (b) simulated magnitude; (c) measured phase; (d) measured magnitude

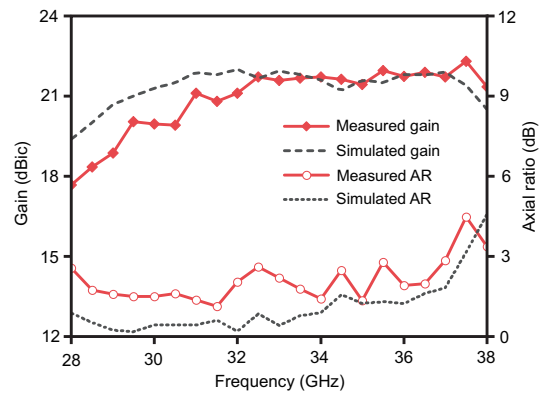


**Fig. 8** Mode purity of the vortex beam versus frequency (OAM: orbital angular momentum)

of different OAM antennas are listed in Table 3. Compared with Bi et al. (2018), Lin et al. (2020), and Wu GB et al. (2020), the proposed TA using sub-wavelength elements achieves more accurate phase compensation and higher mode purity. Although the size of elements used in Zhang FH et al. (2019) is smaller, the proposed element arranged in the honeycomb shape shows more stable performance at the edge of the operating band. For this reason, the proposed TA has a wider bandwidth and higher mode purity compared with the antennas in Zhang FH et al. (2019).

### 5 Conclusions

This correspondence describes a wideband element and a TA for generating a high-purity vortex beam. As validated by fabrication and measurement, the radiated beam maintains typical vortex characteristics in a wide band from 28.5 to 38 GHz. Specifically, OAM mode purity is  $>0.93$  within 28.4% mode purity bandwidth. Due to the wideband and high purity properties, the proposed TA antenna will find several applications, such as large-capacity wireless communication and high-quality radar imaging.



**Fig. 9** Simulation and measurement results of gain and axial ratio (AR)

**Table 3** Comparison with other reported orbital angular momentum (OAM) antennas

Reference	Type	$f_0$ (GHz)	Period of element ( $\times \lambda_0$ )	Aperture ( $\times \lambda_0^2$ )	Mode purity	Mode purity bandwidth (%)*
Zhang FH et al. (2019)	TA	10	0.33	$5 \times 5$	$>0.87$	20.0
Bi et al. (2018)	RA	63	0.63	$31 \times 31$	$>0.79$	17.1
Lin et al. (2020)	RA	23	0.77	$15 \times 15$	$>0.75$	43.4
Wu GB et al. (2020)	TA	300	0.50	$13 \times 13$	$>0.40$	31.3
<b>This work</b>	<b>TA</b>	<b>33.5</b>	<b>0.40</b>	<b><math>10 \times 10</math></b>	<b><math>&gt;0.93</math></b>	<b>28.4</b>

\* Mode purity bandwidth refers to the frequency range in which the OAM mode purity is greater than the value of the previous column. TA: transmitarray; RA: reflectarray

## Contributors

Liangjie QIU designed the research and performed the simulations. Liangjie QIU, Zihang QI, and Yuhan HUANG processed the data. Liangjie QIU drafted the paper. Xiuping LI and Wenyu ZHAO helped organize the paper. Liangjie QIU, Xiuping LI, Zihang QI, and Wenyu ZHAO revised and finalized the paper.

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## Compliance with ethics guidelines

Liangjie QIU, Xiuping LI, Zihang QI, Wenyu ZHAO, and Yuhan HUANG declare that they have no conflict of interest.

## Data availability

The data that support the findings of this study are available from the corresponding author upon reasonable request.

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